

# Entropy-entropy dissipation inequalities for fragmentation-drift equations

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# Fragmentation

Consider a set of particles of different sizes that can break into several pieces, in some random fashion.

$f(t, x) \equiv f(x) \equiv$  density of clusters of size  $x > 0$  at time  $t \geq 0$ .

$b(x, y) \equiv$  probability of obtaining a cluster of size  $y$  when a cluster of size  $x$  breaks.

## Fragmentation equation

$$\partial_t f(x) = Ff(x) := \int_x^\infty b(y, x) f(y) dy - f(x) \int_0^x \frac{y}{x} b(x, y) dy$$

## Typical coefficients

The coefficient of the negative term is

$$B(x) := \int_0^x \frac{y}{x} b(x, y) dy,$$

the *total fragmentation rate*. It is the one for which fragmentation conserves the *total mass*:

$$\int_0^\infty x Ff(x) dx = 0.$$

Typical fragmentation coefficient: think of...

$$b(x, y) = x^{\gamma-1} \quad \text{for some } \gamma \geq 0, \text{ or}$$
$$b(x, y) = \delta_{y=x/2} \quad (\textit{mitosis case})$$

## Self-similar fragmentation

Under fragmentation, the size of clusters becomes smaller and smaller. We rescale to study the long-time shape of the distribution, as it concentrates around size 0:

$$n(t, x) := e^{-2t} f(e^{\gamma t} - 1, e^{-t} x) \quad (t, x > 0).$$

Then,

Self-similar fragmentation equation

$$\partial_t n + x \partial_x n + 2n = \gamma F n$$

# Fragmentation-growth

Now, consider a set of particles of different sizes (cells?) that can randomly break into several pieces, and that also grow at a fixed rate.

$f(t, x), b(x, y)$  as before

Fragmentation-growth equation

$$\partial_t f = -\partial_x f + Ff$$

## Rescaled growth-fragmentation

Distribution grows large in time: we rescale as  $n(t, x) := e^{-\lambda t} f(t, x)$  to get

Rescaled fragmentation-growth equation

$$\partial_t n = -\lambda n - \partial_x n + F n$$

The precise  $\lambda$  must be chosen for this equation to have equilibria: take  $\lambda$  to be the largest eigenvalue of the operator in the r.h.s. (Malthusian growth rate of the population)

## The dual equation: eigenproblem

Find  $\lambda, \phi$  such that:

$$\begin{aligned} -\partial_x \phi &= \int_0^x b(x, y) \phi(y) dy - (B(x) + \lambda) \phi(x), \\ \phi &\geq 0, \quad \int_0^\infty N(x) \phi(x) dx = 1. \end{aligned}$$

Estimates can be obtained through a maximum/minimum principle.

# The eigenproblem

In the case of fragmentation, the dual equation is very easy:

$$\phi(x) = x.$$

In the case of growth-fragmentation, the dual equation is in general difficult to solve: no evident conserved quantity.

- The case of a constant total fragmentation rate is the simplest:  $\lambda = 2$ ,  $\phi(x) \equiv 1$ .
- Otherwise, one has to adjust the  $\lambda$  to the right value and solve for  $\phi$ , a problem studied by Michel (2006), Doumic & Gabriel (2010). With careful estimates, we are able to show that  $\phi(x)$  is bounded below, and looks like  $x$  for  $x$  large.

## Asymptotic behavior

So once rescaled, we look at either

$$\partial_t n = -2n - x \partial_x n + \gamma F n$$

or

$$\partial_t n = -\lambda n - \partial_x n + F n$$

As  $t \rightarrow +\infty$ ,  $n$  converges to a limiting shape (self-similar profile).

*How fast?*

or:

*Do the linear operators on the right hand sides have spectral gaps?*

## Existing results: fragmentation

- Existence and uniqueness of self-similar profiles, and some of their properties (behavior near  $x = 0$  and  $x = +\infty$ ) was proved by Escobedo, Mischler and Rodríguez Ricard (2004/2005).
- Convergence to equilibrium was also proved here using entropy methods (more later).
- No results on rate.

## Existing results: growth-fragmentation

- Perthame and Ryzhik (2005) proved the above for the mitosis case, with constant total fragmentation rate, or perturbations of these. They showed an exponential rate of convergence for these coefficients.
- Michel (2006) and Doumic & Gabriel (2010) studied, for general coefficients, the dual eigenproblem for the growth-fragmentation equation, which essentially implies convergence to equilibrium via entropy methods.
- Laurençot and Perthame (2009) proved exponential decay for a constant total fragmentation rate, allowing for more general kinds of breakage than mitosis.

## Entropy method

In a quite general setting, linear equations which come from a Markov process have many entropy functionals: in our case, take the quadratic one:

$$\begin{aligned} \frac{d}{dt} \int_0^\infty \left( \frac{n(x)}{N(x)} - 1 \right)^2 N(x) \phi(x) dx = \\ - \int_0^\infty \int_y^\infty \phi(y) b(x, y) N(x) \left( \frac{n(x)}{N(x)} - \frac{n(y)}{N(y)} \right)^2 dx dy \end{aligned}$$

$\phi \equiv$  a function obtained from the solution of an eigenvalue problem (more to follow),

$N \equiv$  equilibrium, normalized so that  $\int N\phi = 1$ .

## General relative entropy

A general way to find entropy functionals in many linear equations was given by Michel, Mischler & Perthame (2004/2005). If you have a linear equation with a differential and a “scattering” part (like ours)

$$\partial_t n = Ln$$

assume you can find a stationary profile  $N$  for our equation, and  $\phi$  of the dual operator  $L^*$ :

$$LN = 0, \quad L^*\phi = 0.$$

Then, for any convex  $H$ , this is an entropy of the equation:

$$\int H\left(\frac{n(x)}{N(x)}\right) N(x)\phi(x) dx.$$

## Entropy method (II)

Call  $H$ ,  $D$  the entropy and dissipation of entropy, resp. Then,

$$\frac{d}{dt}H = -D,$$

and if we prove the following functional inequality

$$H \leq CD$$

for some constant  $C$ , we have a speed of convergence to equilibrium:

$$H(t) \leq e^{-Ct}H(0).$$

(Observe that  $\sqrt{H} = \|n - N\|_2$ , where the  $L^2$  norm is w.r.t. the measure  $\phi(x)/N(x) dx$ )

# Functional inequality

So we are led to look at the following inequality:

## Entropy-entropy dissipation inequality

$$\int_0^\infty (u(x) - 1)^2 N(x) \phi(x) dx \leq C \int_0^\infty \int_y^\infty \phi(y) b(x, y) N(x) (u(x) - u(y))^2 dx dy,$$

valid for functions  $u$  with  $\int u N \phi = 1$ .

Can we prove it, for some  $C > 0$ ?

**Not always:** there are examples of  $b$  for which  $D = 0$  does not even imply  $H = 0$  (e.g., mitosis:  $b(x, y) = 2\delta_{y=x/2}$ ).

What happens with  $b(x, y) = x^{\gamma-1}$ ?

# A general type of inequality

Step away a bit from the problem: we want inequalities between the (quadratic) entropy and its dissipation **for a Markov process**.

Let us see some examples.

## Some examples

- 1 The kind of inequality we want:

$$\int_x (u(x) - 1)^2 N(x) \leq C \int_x \int_y N(x) K(x, y) (u(x) - u(y))^2$$

- 2 Fokker-Planck:

$$\int_{\mathbb{R}^d} (u(x) - 1)^2 N(x) dx \leq C \int_{\mathbb{R}^d} N(x) |\nabla u(x)|^2$$

- 3 Discrete Markov process:

$$\sum_i (u_i - 1)^2 N_i \leq C \sum_{i,j} N_i K_{i,j} (u_i - u_j)^2$$

## A general strategy

First, “double the variables”:

$$\sum_i N_i (u_i - 1)^2 = \frac{1}{2} \sum_{i,j} N_i N_j (u_i - u_j)^2.$$

(The normalization is  $\sum u_i N_i = 1$ ). Now it looks more similar to the dissipation:

$$\sum_{i,j} N_i K_{i,j} (u_i - u_j)^2$$

- If we can compare the integrands, it works!
- Otherwise, the inequality is trickier, or may be false.

## A general strategy (II)

An idea from Diaconis & Stroock (Ann. Appl. Prob. '91), in the discrete case: for any  $(i, j)$ , choose a path of jumps from  $i$  to  $j$ , with jumps  $(i, i_1), (i_1, i_2), \dots, (i_m, j)$ . Call it  $\gamma_{ij}$ .

$$\begin{aligned} \sum_{i,j} N_i N_j (u_i - u_j)^2 &= \sum_{i,j} N_i N_j \left( \sum_{(k,l) \in \gamma_{ij}} (u_k - u_l) \right)^2 \\ &\leq \sum_{i,j} N_i N_j \|\gamma_{ij}\| \sum_{(k,l) \in \gamma_{ij}} N_k K_{k,l} (u_k - u_l)^2 \\ &= \sum_{k,l} N_k K_{k,l} (u_k - u_l)^2 \sum_{\gamma_{ij} \ni (k,l)} N_i N_j \|\gamma_{ij}\|, \end{aligned}$$

where the “length” of a path  $\gamma_{ij}$  is  $\|\gamma_{ij}\| := \sum_{(k,l) \in \gamma_{ij}} (N_k K_{k,l})^{-1}$ .

## A general strategy: remarks

- Making a good choice of paths is particular to each problem!
- The strategy for the discrete case suggests how to treat continuous ones.
- Finding a constant in the inequality for a finite system which does not depend on the size of the system amounts to showing inequalities in continuous cases.
- An additional possibility: instead of choosing *one* path for each pair, choose *a set of paths* for each pair, and take the mean in the sum above. This may improve the result! (See later).
- A similar strategy is used in the work of Baranger & Mouhot on explicit spectral gaps for the linearized Boltzmann equation.

## An illustration: weighted Poincaré inequalities

Same idea, in a continuous setting: for any two points on  $\mathbb{R}^d$ , choose a continuous path that joins them; e.g., a segment:

$$\begin{aligned} & \int_{\mathbb{R}^d} \int_{\mathbb{R}^d} N(x)N(y)(u(x) - u(y))^2 dx dy \\ &= \int_{\mathbb{R}^d} \int_{\mathbb{R}^d} N(x)N(y) \left( \int_y^x \nabla u(z) dz \right)^2 dx \\ &\leq \int_{\mathbb{R}^d} \int_{\mathbb{R}^d} N(x)N(y)|x - y| \int_y^x |\nabla u(z)|^2 dz dx \\ &= \int_{\mathbb{R}^d} |\nabla u(z)|^2 \int_{[x,y] \ni z} N(x)N(y)|x - y| d(x,y) dz. \end{aligned}$$

## Back to fragmentation

For fragmentation-drift, we wanted to show that:

$$\begin{aligned} & \int_0^\infty \int_y^\infty \phi(y)\phi(x)N(y)N(x) (u(x) - u(y))^2 dx dy \\ & \leq C \int_0^\infty \int_y^\infty \phi(y)b(x,y)N(x) (u(x) - u(y))^2 dx dy. \end{aligned}$$

Here,  $\phi(y)$  is approximately equal to  $y$  (exactly equal for self-similar fragmentation),  $N$  is a nice function with exponential decay at  $+\infty$ , and the tentative inequality should be

$$\begin{aligned} & \int_0^\infty \int_y^\infty xy N(y)N(x) (u(x) - u(y))^2 dx dy \\ & \leq C \int_0^\infty \int_y^\infty y x^{\gamma-1} N(x) (u(x) - u(y))^2 dx dy. \end{aligned}$$

**Can we compare the integrands in our model cases?** So, can we compare...

$$x y N(y)N(x) \leq C y x^{\gamma-1} N(x) \quad (x, y > 0)$$

this is,

$$N(y) \leq C x^{\gamma-2} \quad (x, y > 0)?$$

This clearly fails for  $x$  large!

## Idea of the proof

Let us try the strategy explained before: we pass through “intermediate reactions” in order to compare terms which do not correspond to the same  $x, y$ :

$$u(x) - u(y) = u(x) - u(z) + u(z) - u(y).$$

We need the additional trick: we average among a range of possible splittings:

$$\begin{aligned} u(x) - u(y) &= \int_2^3 (u(x) - u(y/s) + u(y/s) - u(y)) ds \\ &= \int_{y/3}^{y/2} \frac{y}{z^2} (u(x) - u(z) + u(z) - u(y)) dz. \end{aligned}$$

Then, by Cauchy-Schwarz's inequality,

$$\begin{aligned}(u(x) - u(y))^2 &\leq \frac{y}{6} \int_{y/3}^{y/2} \frac{y^2}{z^4} (u(x) - u(z) + u(z) - u(y))^2 dz \\ &\leq \frac{27}{2y} \int_{y/3}^{y/2} (u(x) - u(z) + u(z) - u(y))^2 dz \\ &\leq \frac{27}{y} \int_{y/3}^{y/2} (u(x) - u(z))^2 dz + \frac{27}{y} \int_{y/3}^{y/2} (u(z) - u(y))^2 dz.\end{aligned}$$

One then has to plug these two terms in the l.h.s. of the inequality, and exchange the order of integration. But the gain is clear here: one can gain up to one order in the comparison we tried before.

## First term

$$\begin{aligned} & \int_0^\infty \int_{\max\{4x, M\}}^\infty x N(x) N(y) \int_{y/3}^{y/2} (u(x) - u(z))^2 dz dy dx \\ &= \int_0^\infty \int_{\max\{\frac{4x}{3}, \frac{M}{3}\}}^\infty x N(x) (u(x) - u(z))^2 \int_{\max\{4x, 2z, M\}}^{3z} N(y) dy dz dx \\ &\leq C \int_0^\infty \int_{\max\{\frac{4x}{3}, \frac{M}{3}\}}^\infty x N(x) (u(x) - u(z))^2 \frac{1}{1+z} \int_{2z}^\infty y N(y) dy dz dx \\ &\leq C \int_0^\infty \int_{\frac{4x}{3}}^\infty x N(x) \frac{N(z)}{1+z} (u(x) - u(z))^2 dz dx \leq C D, \end{aligned}$$

## Second term

$$\begin{aligned} & \int_0^\infty \int_{\max\{4x, M\}}^\infty x N(x)N(y) \int_{y/3}^{y/2} (u(z) - u(y))^2 dz dy dx \\ & \leq C \int_0^\infty \int_{4x}^\infty x N(x)N(y) \int_{y/3}^{y/2} (u(z) - u(y))^2 dz dy dx \\ & = C \int_0^\infty \int_{y/3}^{y/2} N(y)(u(z) - u(y))^2 \int_0^{y/4} xN(x) dx dz dy \\ & \leq C \int_0^\infty \int_{y/3}^{y/2} N(y)(u(z) - u(y))^2 dz dy \\ & \leq C \int_0^\infty \int_{y/3}^{y/2} \frac{z}{1+y} N(y)(u(z) - u(y))^2 dz dy \leq CD. \end{aligned}$$

## Theorem

*The entropy-entropy dissipation inequality mentioned before is true for  $b(x, y) = y^{\gamma-1}$  for the range of  $\gamma$  below.*

*Hence: assume*

- $0 \leq \gamma < 2$  for growth-fragmentation.
- $0 < \gamma < 2$  for self-similar fragmentation.

*Then, for any solution  $n$  to one of these models,*

$$\|n - N\|_2 \leq e^{-Ct} \|n_0 - N\|_2,$$

*for some constant  $C > 0$ .*

## Estimates we need

For the previous proof to work for a given stationary profile  $N$  and dual profile  $\phi$  we need:

- 1 Exponential decay estimates on  $N$  (by above and below).
- 2 Accurate estimates on  $\phi$  (above and below by some power).

Let us do them in the growth-fragmentation case (more difficult for this).

## Estimates on $N$ : moments

First, study the moments of  $N$ : if  $M_k$  is the moment of order  $k$ ,

$$\left(1 - \frac{1}{k}\right)M_{\gamma+k} \leq kM_{k-1}.$$

Summing to get exponential moments,

$$\int_0^\infty \exp(ax^{\gamma+1})N(x) dx < +\infty.$$

## Estimates on $N$ : bounds above and below

Trick: define

$$\Lambda(x) := \lambda x + \int_0^x B(x) dx.$$

When  $b(x, y) = 2x^{\gamma-1}$ , then  $N(x) = e^{-\Lambda(x)}$  is *an exact solution*.

Otherwise, we calculate:

$$(Ge^{a\Lambda})' = (a-1)(B + \lambda)Ge^{a\Lambda} + e^{a\Lambda} \mathcal{L}_+(G).$$

- Right hand side integrable  $\implies$  upper bound.
- Right hand side nonnegative  $\implies$  lower bound.

## Estimates on $N$ : result

Finally, this gives

$$C_1 e^{-a_1 x^{\gamma+1}} \leq G(x) \leq C_2 e^{-a_2 x^{\gamma+1}} \quad (x > \delta)$$

for some constants  $a_1, a_2, C_1, C_2 > 0$ .

## Maximum principle for $\phi$

$\phi$  is limit of solutions of

$$0 = S\phi_L(x) := -\phi_L'(x) + (\lambda_L + B(x))\phi_L(x) - \int_0^x b(x, y)\phi_L(y) dy.$$

Lemma (Maximum principle for  $S$ )

Take  $A \geq 1/\lambda_L$ . If  $Sw \geq 0$  on  $(A, L)$ ,  $w \geq 0$  on  $[0, A]$  and  $w(L) \geq 0$  then  $w \geq 0$  on  $[A, L]$ .

## Maximum principle: proof

Assume  $x_0 \in (A, L)$  with  $w(x_0) < 0$ , and such that  $w(x_0)/x_0$  is minimum. Then  $w'(x_0) = w(x_0)/x_0$ , so

$$\begin{aligned} \mathcal{S}(x_0) &= -\frac{w(x_0)}{x_0} + (\lambda_L + B(x_0)) w(x_0) - \int_0^{x_0} b(x_0, y) w(y) dy \\ &\leq -\frac{w(x_0)}{x_0} + (\lambda_L + B(x_0)) w(x_0) - \frac{w(x_0)}{x_0} \int_0^{x_0} b(x_0, y) y dy \\ &= w(x_0) \left( \lambda_L - \frac{1}{x_0} \right) < 0, \end{aligned}$$

which contradicts that  $\mathcal{S}w \geq 0$  on  $(A, L)$ .

## Estimates for $\phi$ (result)

By comparing with  $w(x) = C(1 + x)$  and  $w(x) = x^k \phi$  (with  $k < 1$  and  $\phi$  a truncation so that  $w(L) = 0$ ), we can prove

$$C_k x^k \leq \phi(x) \leq C(1 + x)$$

when  $\gamma > 0$ .

Thanks!