

AdS Solutions Through Transgression

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Supersymmetric AdS/CFT Solutions

Programme: characterise general classes of supersymmetric warped product solutions of D=11 or type IIB supergravity of the form:

$AdS_{d+1} \times_w M$:

$$ds^2 = e^{2A(y)} [ds^2(AdS_{d+1}) + ds^2(M)(y)]$$

with fluxes preserving isometries of AdS .

Solutions are invariant under $SO(d, 2) \Leftrightarrow$ characterise the most general SCFTs in d spacetimes that have a supergravity description.

Motivation

- ★ SCFTs are interesting for many reasons.
- ★ AdS_3 , AdS_2 solutions \leftrightarrow blackholes.
- ★ Can analytically continue \rightarrow “BPS bubbles” dual to smooth BPS states of SCFTs (LLM) Also: Wilson Lines Lunin, D'Hoker, Gutperle

Programme

1. Use G -Structure techniques JPG, Martelli, Pakis, Waldram (“Invariant Tensor Method”, “Spinorial Geometry”, ...) to give a precise characterisation of geometry ($ds^2(M)$, A , fluxes).
Prove Theorems.
2. In favourable cases construct explicit solutions and try to identify the SCFTs.

e.g. $AdS_5 \times SE_5$ solutions in type IIB and N=1 SCFTs.

Special Cases

AdS_3 in Type IIB sugra:

$$ds^2 = e^{-B/2}[ds^2(AdS_3) + ds^2(Y_7)]$$

$$F_5 = (1 + *)Vol(AdS_3) \wedge F_2$$

Dual to $N=(0,2)$ SCFTs in $D=2$: D3-branes on 2-cycles in CY_4 .

AdS_2 in $D=11$ sugra:

$$ds^2 = e^{-2B/3}[ds^2(AdS_2) + ds^2(Y_9)]$$

$$G_4 = Vol(AdS_2) \wedge F_2$$

Dual to $N=2$ SCQM: M2-branes wrapped on 2-cycles in CY_5 .

Geometry given by $ds^2(Y_{2n+1}), F_2, B$.

Also: 1/8 BPS Bubbles with $SO(4)$ or $SO(3)$ symmetry that generalise LLM .

There is always a Killing vector - dual to R symmetry
Locally, metrics can be written Kim

$$ds_{2n+1}^2 = \frac{(n-2)^2}{4}(dz + P)^2 + e^B ds^2(B_{2n})$$

where B_{2n} is Kähler and satisfies

$$\square R - \frac{1}{2}R^2 + R_{ij}R^{ij} = 0 \quad (*)$$

which arises from the Bianchi identity for the five-form $dF_5 = 0$.

Furthermore, e^B and F_2 are fixed by Kähler geometry and $dP = \mathcal{R}$.

Many similarities to $AdS_5 \times SE_5$ solutions of type IIB and $AdS_4 \times SE_7$ solutions of D=11.

$$ds^2(SE_{2n+1}) = (dz + P)^2 + ds^2(B_{2n})$$

where B_{2n} is Kähler and Einstein and $dP = \mathcal{R}$.

Can we switch on additional fluxes? Recall Klebanov and Strassler: start with $AdS_5 \times T^{1,1}$ solution of type IIB. Observe that the cone metric C over $T^{1,1}$ is the Calabi-Yau conifold metric. Then observe that there is a $\mathbb{R}^{1,3} \times_w C$ solution with $G \neq 0$ provided that $dG = 0$, G is (1,2) and $J \wedge G = 0$. Solution boils down to solving an equation for a single function that arises from the Bianchi identity with “transgression terms”

$$dF_5 = \frac{i}{2} G \wedge G^*$$

” Brane Resolution Through Transgression” [Cvetic, Lu, Pope](#)

AdS_3 in Type IIB sugra:

$$ds^2 = e^{-B/2}[ds^2(AdS_3) + ds^2(Y_7)]$$

$$F_5 = (1 + *)Vol(AdS_3) \wedge F_2$$

$$G \neq 0$$

Dual to $N=(0,2)$ SCFTs in $D=2$:

$$ds_7^2 = \frac{1}{4}(dz + P)^2 + e^B ds^2(B_6)$$

Require G to be a closed, primitive, $(1, 2)$ form on B_6 . If such a form exists the entire solution boils down to solving

$$\square R - \frac{1}{2}R^2 + R_{ij}R^{ij} + \frac{2}{3}G^{ijk}G_{ijk}^* = 0 \quad (*)$$

with $dP = \mathcal{R}$. which arises from solving the modified Bianchi identity $dF_5 = \frac{i}{2}G \wedge G^*$.

AdS_2 in D=11 sugra:

$$ds^2 = e^{-2B/3} [ds^2(AdS_3) + ds^2(Y_7)]$$

$$G_4 = Vol(AdS_2) \wedge F_2 + F_4$$

Dual to N=(0,2) SCFTs in D=2:

$$ds_9^2 = (dz + P)^2 + e^B ds^2(B_8)$$

Require F_4 to be a closed, primitive, (2,2) form on B_8 . If such a form exists the entire solution boils down to solving

$$\square R - \frac{1}{2}R^2 + R_{ij}R^{ij} + \frac{1}{4!}F_4^{ijkl}F_{4ijkl} = 0 \quad (*)$$

with $dP = \mathcal{R}$. which arises from solving the modified Bianchi identity $dG_4 = \frac{1}{2}G_4 \wedge G_4$.

Explicit Solutions

Adapt constructions of Sasaki-Einstein metrics $Y^{p,q}$ [JPG, Martelli, Sparks, Waldram](#)

Use [Page, Pope](#) to build Kähler bases by considering line bundles over Kähler-Einstein manifolds and then try and solve (*).

In more detail

$$ds_{2n}^2 = \frac{1}{U} d\rho^2 + U \rho^2 D\phi^2 + \rho^2 ds^2(KE_{2n-2})$$

where $(D\phi) = 2J_{KE}$. Demanding that it solves (*) leads to solving a second order non-linear ODE for U . Remarkably it admits polynomial solutions!

Type IIB: $G = 0$

Infinite classes of solutions when we use ds_6^2 built from $KE_4 = CP^2$, $KE_4 = S^2 \times S^2$ or $B_4 = S^2 \times H^2$.

Type IIB: $G \neq 0$

Let $ds_6^2 = ds_4^2 \times T^2$ where ds_4^2 is built from a line bundle over S^2 .
Then set

$$G = d\bar{u} \wedge d\left(\frac{q}{\rho^2} D\phi\right)$$

Leads to infinite new solutions parametrised by q .

For these solutions with a T^2 factor we can T-dualise on one leg and then on the other leg of the T^2 to get new type IIB solutions. After a further S -duality we find that the solutions are constructed from NS fields only!

$$\begin{aligned} ds^2 &= ds^2(AdS_3) \times_w ds^2(X_7) \\ e^{2\phi} &\neq 0 \\ H_3 &\neq 0 \end{aligned}$$

Furthermore $X^7 = S^3 \times S^3 \times S^1$.

Final Comments

- ★ Studied general class of AdS_3 solutions of type IIB and AdS_2 solutions of D=11 supergravity.
- ★ Many similarities with Sasakian Einstein Geometry.
- ★ Motivated generalisations with $G \neq 0$ and magnetic F_4 .
- ★ Motivated construction of explicit solutions - infinite new examples both for type IIB and D=11.
- ★ For type IIB there are duality frames where some of the solutions are constructed from NS field only: perhaps we can make progress on constructing the bulk CFT?

★ What are dual $d = 2$ (0,2) SCFTs? Are they related to SCFTs that are dual to SE ? In some cases, it appears that switching on the three-form corresponds to exactly marginal deformations (there are also additional beta deformations).

★ Bubble solutions.

★ Extensive classification of supersymmetric AdS_n geometries, for various n in type IIB and $D = 11$. but more to be done. For some classes there are analogously rich classes of explicit solutions - to do with the fact that the solutions are complex?