

Viscosity Solutions for Nonlinear Hyperbolic Systems

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Hyperbolic Systems of Conservation Laws

$$u_t + f(u)_x = 0$$

$u = (u_1, \dots, u_n) \in \mathbb{R}^n$ conserved quantities

$f = (f_1, \dots, f_n) : \mathbb{R}^n \mapsto \mathbb{R}^n$ fluxes

For gas-dynamics, conserved quantities are:

mass momentum energy

$$u_t + A(u)u_x = 0 \qquad A(u) = Df(u)$$

STRICTLY HYPERBOLIC: each matrix $A(u)$ has real distinct eigenvalues

$$\lambda_1(u) < \lambda_2(u) < \dots < \lambda_n(u)$$

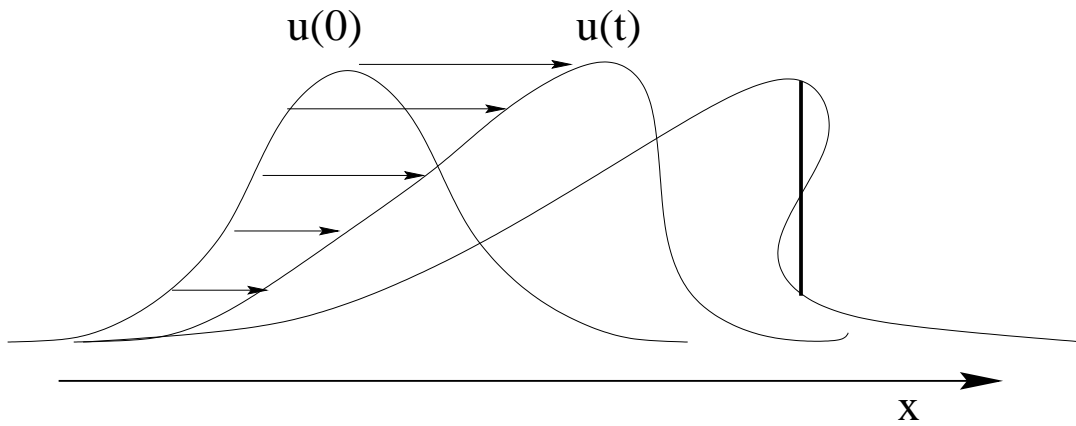
eigenvectors $r_1(u), \dots, r_n(u)$

GENUINELY NONLINEAR: $D\lambda_i \cdot r_i > 0$

LINEARLY DEGENERATE: $D\lambda_i \cdot r_i \equiv 0$

Loss of Regularity

$$u_t + f(u)_x = 0 \quad u(0, x) = \bar{u}(x)$$



Global solutions: in a space of discontinuous functions

$$u(t, \cdot) \in BV \text{ or } \mathbf{L}_{loc}^1$$

Conservation equations interpreted in a distributional sense

$$\iint \{u\phi_t + f(u)\phi_x\} dxdt = 0 \quad \phi \in \mathcal{C}_c^1$$

Construction of Solutions

Standard techniques do not apply:

- Fixed point theorems
- Variational principles
- Upper, lower solutions

Two approaches:

The BV setting (J.Glimm, 1965)

- Approximate solutions are constructed in a class of piecewise constant functions
- Total variation is controlled by an Interaction Potential

The L^p setting (L.Tartar, R.DiPerna, 1979-1983)

- Consider vanishing viscosity approximations

$$u_t^\varepsilon + f(u^\varepsilon)_x = \varepsilon u_{xx}^\varepsilon$$

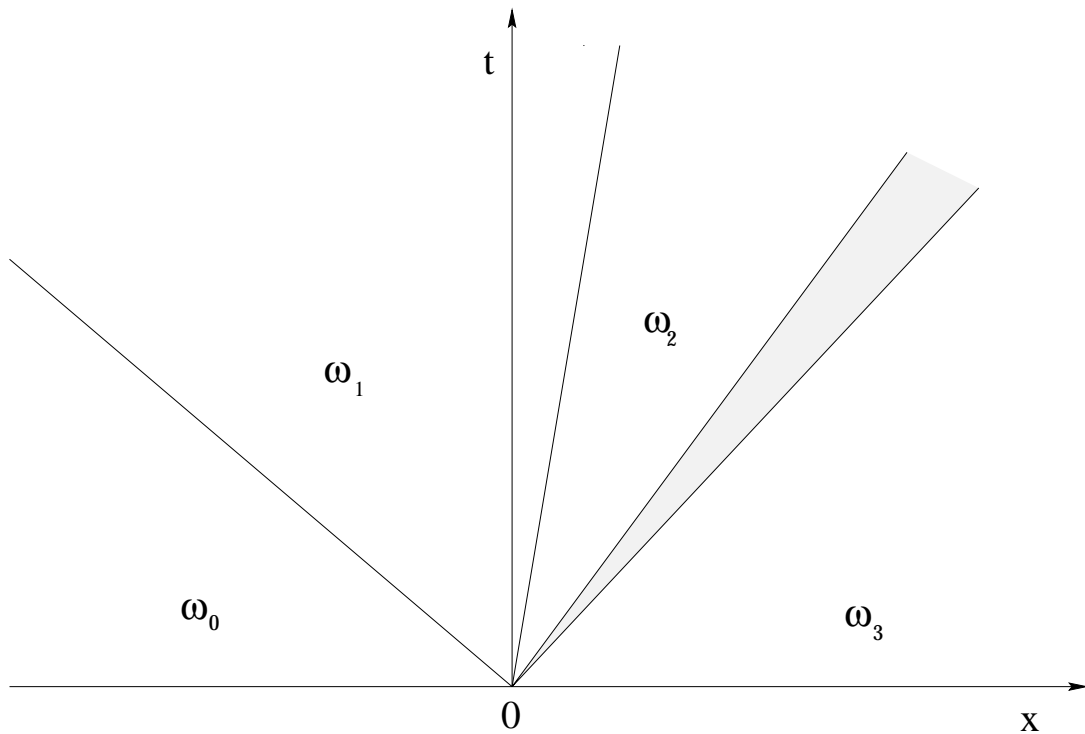
- A bound in L^p yields a weakly converging sequence $u^\varepsilon \rightharpoonup u$ as $\varepsilon \rightarrow 0$
- By compensated compactness argument, u is a weak solution of the nonlinear system of conservation laws.
- Mainly restricted to 2×2 systems

Construction of BV solutions

Building block: the Riemann Problem

$$u_t + f(u)_x = 0$$
$$u(0, x) = \begin{cases} u^- & \text{if } x < 0 \\ u^+ & \text{if } x > 0 \end{cases}$$

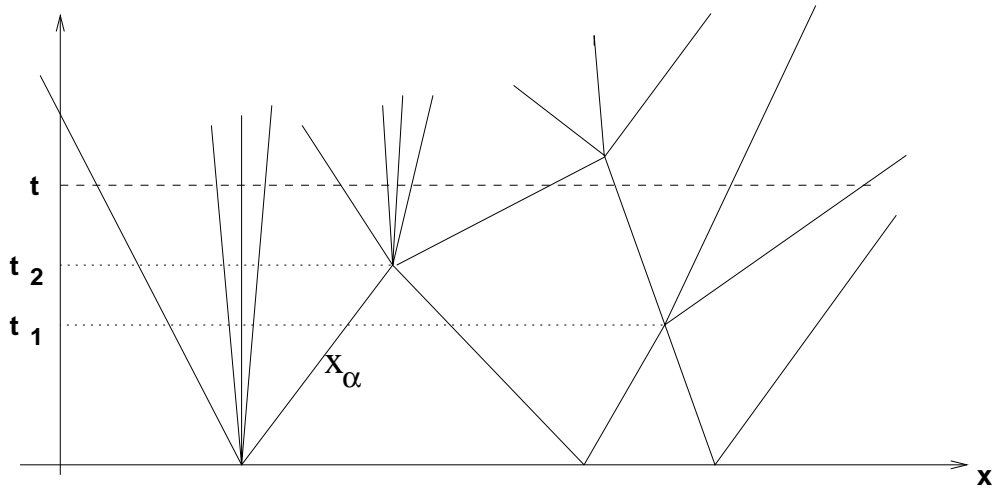
Self-similar solution (P. Lax, 1957)



The Cauchy problem

$$u_t + f(u)_x = 0, \quad u(0, x) = \bar{u}(x)$$

Front Tracking Approximations



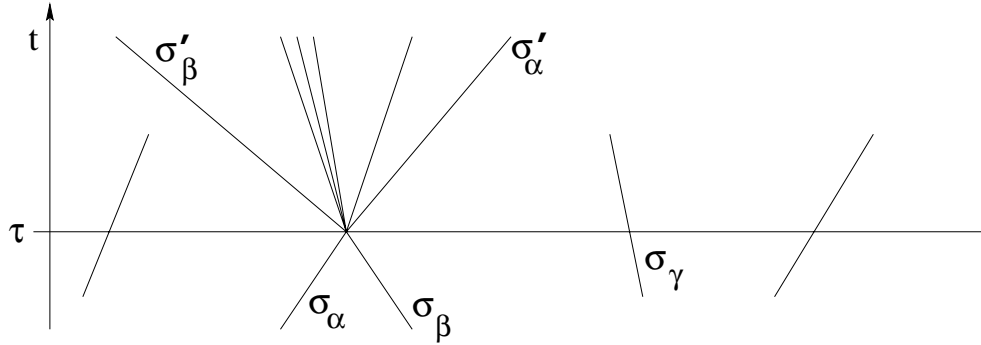
piecewise constant ε -approximate solution

$$\sum_{\alpha} \left| \dot{x}_{\alpha} (u_{\alpha}^{+} - u_{\alpha}^{-}) - (f(u_{\alpha}^{+}) - f(u_{\alpha}^{-})) \right| < \varepsilon \quad \text{for a.e. } t > 0$$

Uniform BV bounds \implies compactness

\implies a subsequence converges in \mathbf{L}_{loc}^1 to a weak solution

Glimm interaction functional $Q(u) = \sum_{\text{approaching}} |\sigma_\alpha \sigma_\beta|$



Uniqueness and Continuous Dependence

$$u_t + f(u)_x = 0 \quad u(0, x) = \bar{u}(x) \quad (1)$$

Theorem 1 (Stability) *The solution of the Cauchy problem (1) obtained as limit of front tracking approximations is unique and depends Lipschitz continuously on the initial data.*

$$\|u(t, \cdot) - v(t, \cdot)\|_{L^1} \leq L \cdot \|\bar{u} - \bar{v}\|_{L^1} \quad t \geq 0$$

2×2
 $n \times n$
 $n \times n$

A.B. - R.M.Colombo (1994)
A.B. - B.Piccoli (1996)
A.B. - T.P.Liu - T.Yang (1998)

Lyapunov Functional

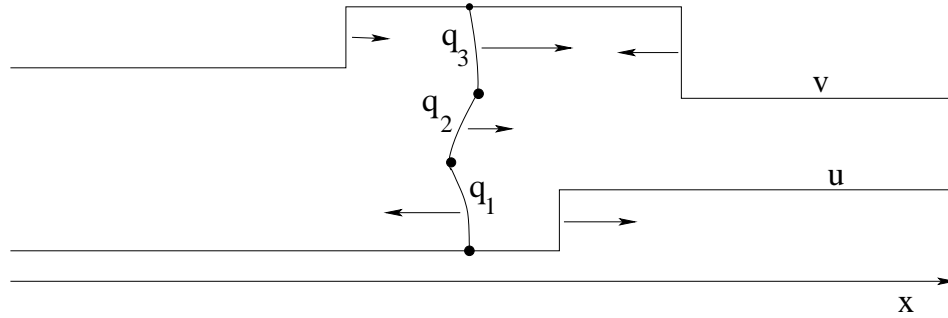
$$u_t + f(u)_x = 0$$

$$\frac{1}{C} \cdot \|v - u\|_{\mathbf{L}^1} \leq \Phi(u, v) \leq C \cdot \|v - u\|_{\mathbf{L}^1}$$

$$\frac{d}{dt} \Phi(u(t), v(t)) \leq 0$$

Construction of $\Phi(u, v)$ when u, v piecewise constant

$q_i(x)$ = i-th component of the jump $v(x) - u(x)$



$$\Phi(u, v) \doteq \sum_{i=1}^n \int_{-\infty}^{\infty} |q_i(x)| W_i(x) dx$$

$$\begin{aligned} W_i(x) &\doteq 1 + \kappa_1 \cdot [\text{total strength of waves in } u \text{ and in } v \text{ which approach } q_i(x)] \\ &\quad + \kappa_2 \cdot [\text{wave interaction potentials of } u \text{ and of } v] \\ &\doteq 1 + \kappa_1 A_i(x) + \kappa_2 [Q(u) + Q(v)] \end{aligned}$$

Uniqueness

$$u_t + f(u)_x = 0 \quad u(0, x) = \bar{u}(x) \quad (1)$$

Semigroup of solutions obtained as limits of front-tracking approximations

$$S : \mathcal{D} \times [0, \infty[\mapsto \mathcal{D} \quad t \mapsto u(t, \cdot) = S_t \bar{u}$$

$$\|S_t \bar{u} - S_t \bar{v}\|_{\mathbf{L}^1} \leq L \cdot \|\bar{u} - \bar{v}\|_{\mathbf{L}^1}$$

Theorem 2 (Uniqueness). *Let the map $u = u(t, x)$ be a weak solution of the Cauchy problem (1), with small total variation, which satisfies*

Lax Entropy Condition: *At every point of approximate jump, one has*

$$\lambda_i(u^+) \leq \lambda \leq \lambda_i(u^-)$$

Bounded Variation Condition: *u has finite total variation along every space-like curve γ in the t - x plane.*

Then $u(t, \cdot) = S_t \bar{u}$ for all $t \in [0, T]$, hence the solution is unique.

(A.B., P.LeFloch, P.Goatin, M.Lewicka, 1997-2000)

A. Bressan, *Hyperbolic Systems of Conservation Laws. The One Dimensional Cauchy Problem*, Oxford University Press 2000.

Vanishing Viscosity Approximations

$$u_t^\varepsilon + f(u^\varepsilon)_x = \varepsilon u_{xx}^\varepsilon \quad u(0, x) = \bar{u}(x) \quad (1)$$

Conjecture: as $\varepsilon \rightarrow 0$, the solution u^ε of (1) converges to the unique entropy weak solution of

$$u_t + f(u)_x = 0 \quad u(0, x) = \bar{u}(x) \quad (2)$$

1. Scalar case (S.Kruzhkov, 1970)

- Based on comparison arguments
- Valid in several space dimensions.

2. Compensated Compactness (L.Tartar - R.DiPerna, 1979-83)

A subsequence u^ε converges (weakly) to some weak solution u of (2).

- applies mainly to 2×2 systems
- $u \in \mathbf{L}^\infty$, possibly $u \notin BV$. Uniqueness, stability ??

3. Singular Perturbations (J.Goodman - Z.Xin, 1992)

If u is piecewise smooth entropic solution of (2), then there exists a convergent sequence $u^\varepsilon \rightarrow u$.

- for $n \times n$ systems, piecewise smooth solutions are not dense

Key problem: derive BV bounds on $u^\varepsilon(t, \cdot)$, independent of ε . This will imply strong convergence $\|u^\varepsilon - u\|_{\mathbf{L}^1_{loc}} \rightarrow 0$ and uniqueness.

A priori BV bounds

The rescaling $t \mapsto t/\varepsilon$, $x \mapsto x/\varepsilon$ transforms

$$u_t + A(u)u_x = \varepsilon u_{xx} \quad u(0, x) = \bar{u}(x)$$

into

$$u_t + A(u)u_x = u_{xx} \quad u(0, x) = \bar{u}(\varepsilon x) \quad (3)$$

Assuming \bar{u} small in BV, estimate: $\text{Tot.Var.}\{u(t, \cdot)\} \quad t > 0$

Basic strategy:

1. Decompose the gradient along a suitable basis of unit vectors

$$u_x = \sum v^i \tilde{r}_i$$

2. Write an evolution equation for the gradient components

$$v_t^i + (\tilde{\lambda}_i v^i)_x - v_{xx}^i = \phi_i$$

3. Show that the source terms ϕ_i are INTEGRABLE, hence

$$\|v^i(t, \cdot)\|_{\mathbf{L}^1} \leq \|v^i(0, \cdot)\|_{\mathbf{L}^1} + \int_0^\infty \int_{\mathbb{R}} |\phi_i(s, x)| dx dt$$

4. Conclude that

$$\text{Tot.Var.}\{u(t, \cdot)\} = \|u_x(t, \cdot)\|_{\mathbf{L}^1} \leq \sum \|v^i(t, \cdot)\|_{\mathbf{L}^1} < \infty$$

Tracing Viscous Waves

$$u_t + A(u)u_x = u_{xx} \tag{3}$$

The decomposition

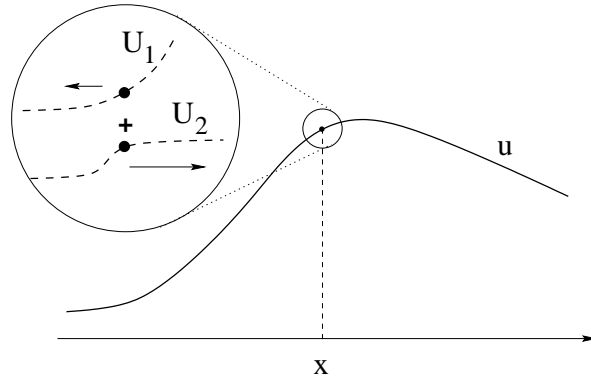
$$v^i = l_i \cdot u_x \qquad u_x = \sum_i v^i r_i \tag{4}$$

along the eigenvectors r_1, \dots, r_n of $A(u)$ yields

$$\begin{aligned} & v_t^i + (\lambda_i v^i)_x - v_{xx}^i \\ &= l_i \cdot \left\{ \sum_{j \neq k} \lambda_j [r_j, r_k] v^j v^k + 2 \sum_{j,k} (r_k \bullet r_j) v_x^j v^k + \sum_{j,k,\ell} [r_\ell, r_k \bullet r_j] v^j v^k v^\ell \right\} \end{aligned}$$

It does NOT work. One should achieve $\phi_i \equiv 0$ in case of a viscous travelling wave.

Problem: Given u, u_x, u_{xx} at a point x , is there a meaningful way to decompose u_x as a sum of gradients of viscous travelling waves ?



STEP 1: Construct a family of travelling wave profiles through a given state $u(x) \in \mathbb{R}^n$, having the right number of parameters to fit the data

$$(u_x, u_{xx}) \in \mathbb{R}^{n+n}$$

$$U_i(x) = u(x) \quad i = 1, \dots, n$$

$$\sum_i U_i'(x) = u_x(x), \quad \sum_i U_i''(x) = u_{xx}(x)$$

For each i , select a 2-parameter family of travelling waves

$$(A(U_i) - \sigma_i)U_i' = U_i''$$

A Center Manifold Approach

(Stefano Bianchini, 2000)

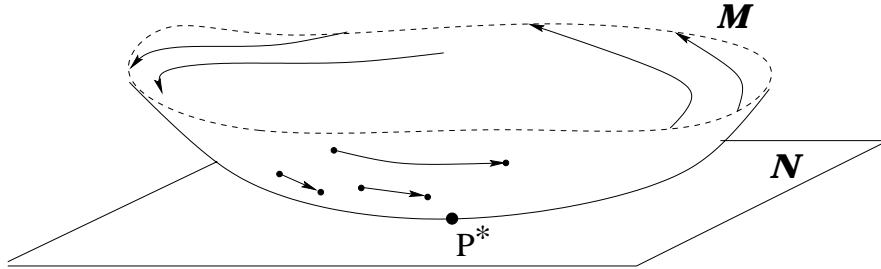
$$u_t + A(u)u_x = u_{xx}$$

travelling wave solution $u(t, x) = U(x - \sigma t)$

$$U'' = (A(U) - \sigma)U' \quad \begin{cases} \dot{u} = v \\ \dot{v} = (A(u) - \sigma)v \\ \dot{\sigma} = 0 \end{cases} \quad (6)$$

Equilibrium point: $P^* \doteq (u^*, 0, \lambda_i(u^*))$

Center Manifold Theorem: In a neighborhood of the point P^* there exists a center manifold \mathcal{M}_i , of dimension $n + 2$, locally invariant for the flow of (6).



For each $(u, v^i, \sigma_i) \approx (u^*, 0, \lambda_i(u^*))$ this construction singles out an i -travelling wave $U(s)$ passing through u (say, at $s = 0$) such that

$$v^i = \pm |U'| = \text{signed strength}$$

$$\sigma_i = \text{speed}$$

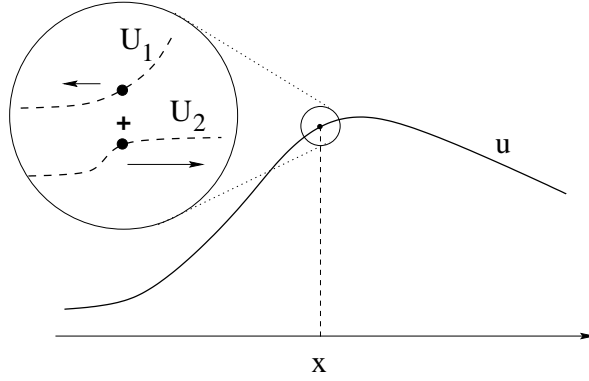
Call $\tilde{r}_i(u, v^i, \sigma_i) = U'/|U'|$ the unit vector tangent to this travelling wave profile

Wave Decomposition

$\tilde{r}_i = \tilde{r}_i(u, v^i, \sigma_i)$ is a smooth map

When $v^i = 0$, one has $\tilde{r}_i(u, 0, \sigma_i) = r_i(u) = i$ -eigenvector of $A(u)$, hence $\{\tilde{r}_1, \dots, \tilde{r}_n\}$ is a basis of linearly independent unit vectors

Given a smooth function $u : \mathbb{R} \mapsto \mathbb{R}^n$, at each point x we seek a decomposition of u_x as a sum of gradients of travelling waves, depending on the second order jet (u, u_x, u_{xx})



Key idea: solve simultaneously the $2n$ equations for $v^1, \dots, v^n, w^1, \dots, w^n$

$$\begin{aligned}
 u_x &= \sum_{i=1}^n v^i \tilde{r}_i(u, v^i, \sigma_i) \\
 u_t \doteq u_{xx} - A(u)u_x &= \sum_{i=1}^n w^i \tilde{r}_i(u, v^i, \sigma_i)
 \end{aligned}
 \quad \sigma_i = -\theta_i \left(\frac{w^i}{v^i} \right) \quad (8)$$

For a j -travelling wave contained in the center manifold, this yields

$$\begin{aligned}
 u_x &= v^j \tilde{r}_j & u_t &= w^j \tilde{r}_j & \sigma_j &= -\frac{w^j}{v^j} = -\frac{u_t}{u_x} \\
 v^i &= 0 & & \text{for } i \neq j & &
 \end{aligned}$$

Evolution of Gradient Components

Given u, u_x, u_{xx} , there exists a unique decomposition

$$u_x = \sum v^i \tilde{r}_i \quad u_t = \sum w^i \tilde{r}_i$$

Defining $\tilde{\lambda}(u, v^i, \sigma_i) \doteq \langle \tilde{r}_i, A(u)\tilde{r}_i \rangle$, after n^n computations one finds

$$\begin{cases} v_t^i + (\tilde{\lambda}_i v^i)_x - v_{xx}^i = \phi_i \\ w_t^i + (\tilde{\lambda}_i w^i)_x - w_{xx}^i = \psi_i \end{cases} \quad (9)$$

$$\begin{aligned} \phi_i, \psi_i = & \mathcal{O}(1) \cdot \sum |w^j + \theta_j v^j| \cdot [|v^j w^j|/\delta_1 + |v_x^j| + |w_x^j|] && \text{(wrong speed)} \\ & + \mathcal{O}(1) \cdot \sum |w_x^j v^j - v_x^j w^j| && \text{(change in strength)} \\ & + \mathcal{O}(1) \cdot \sum \frac{1}{\delta_1} \left| v^j \left(\frac{w^j}{v^j} \right)_x \right|^2 && \text{(change in speed)} \\ & + \mathcal{O}(1) \cdot \sum_{j \neq k} [|v^j v^k| + |v_x^j v^k| + |v^j w^k| + |v_x^j w^k| + |w^j w^k|] \\ & && \text{(interaction of waves of different families)} \end{aligned}$$

KEY ESTIMATES

Wrong speed \implies Energy estimates for the heat equation

Change in wave strength \implies Area functional

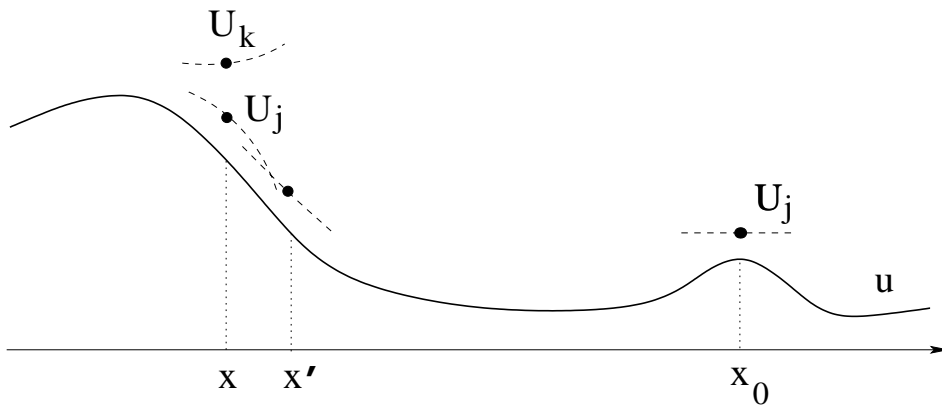
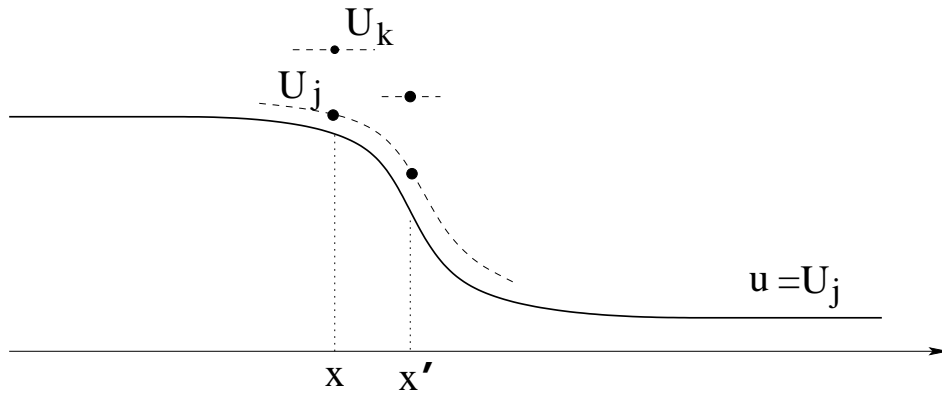
Change in wave speed \implies Curve shortening

Interaction of waves of different families \implies Wave interaction potential

(A.B. and S.Bianchini, 1999-2000)

$$\|v^i(t, \cdot)\|_{\mathbf{L}^1}, \|w^i(t, \cdot)\|_{\mathbf{L}^1}, \|u_x(t, \cdot)\|_{\mathbf{L}^1} = \mathcal{O}(1) \quad \text{for all } t \geq 0$$

Motivation of the Source Terms



Wrong speed: The cut-off function is active, near a point x_0 where $u_x \approx 0$, so that $|u_{xx}| \gg |u_x|$

Change in wave speed, or strength: The viscous travelling j -wave that best approximates u near a point x is not the same at a nearby point x'

Transversal wave interactions: At a point x , waves of distinct families $j \neq k$ are present

Interaction of Viscous Waves of Distinct Families

$$\begin{cases} z_t + [\lambda(t, x)z]_x - z_{xx} = 0 \\ z_t^* + [\lambda^*(t, x)z^*]_x - z_{xx}^* = 0 \end{cases}$$

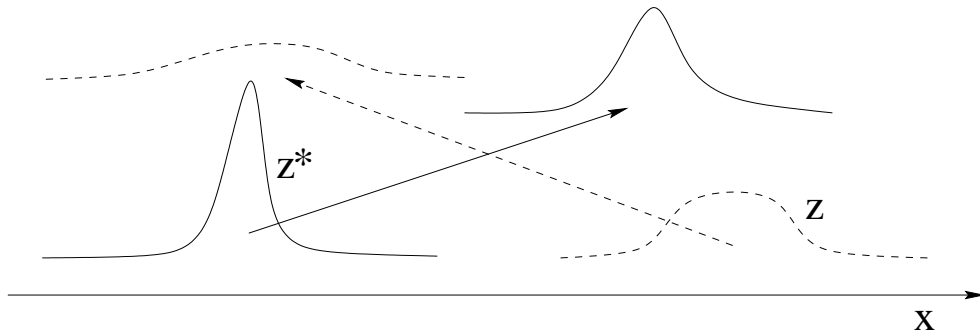
$$\inf_{t,x} \lambda^*(t, x) - \sup_{t,x} \lambda(t, x) \geq c$$

$z \doteq$ density of waves with speed λ

$z^* \doteq$ density of waves with speed λ^*

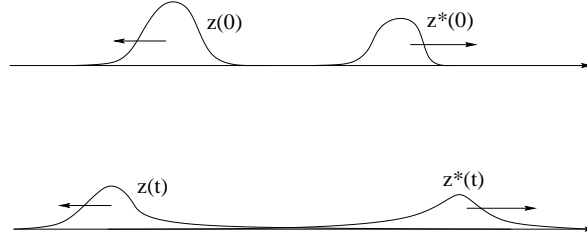
Total amount of interaction:

$$\int_0^\infty \int_{\mathbb{R}} |z(t, x)| \cdot |z^*(t, x)| dx dt$$

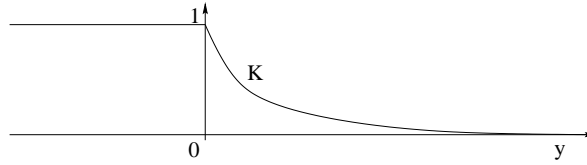


A Potential for Transversal Wave Interactions

$$Q(z, z^*) \doteq \iint_{\mathbb{R}^2} K(x_2 - x_1) |z(x_1)| |z^*(x_2)| dx_1 dx_2$$



$$K(y) \doteq \begin{cases} e^{-cy/2} & \text{if } y > 0 \\ 1 & \text{if } y \leq 0 \end{cases}$$



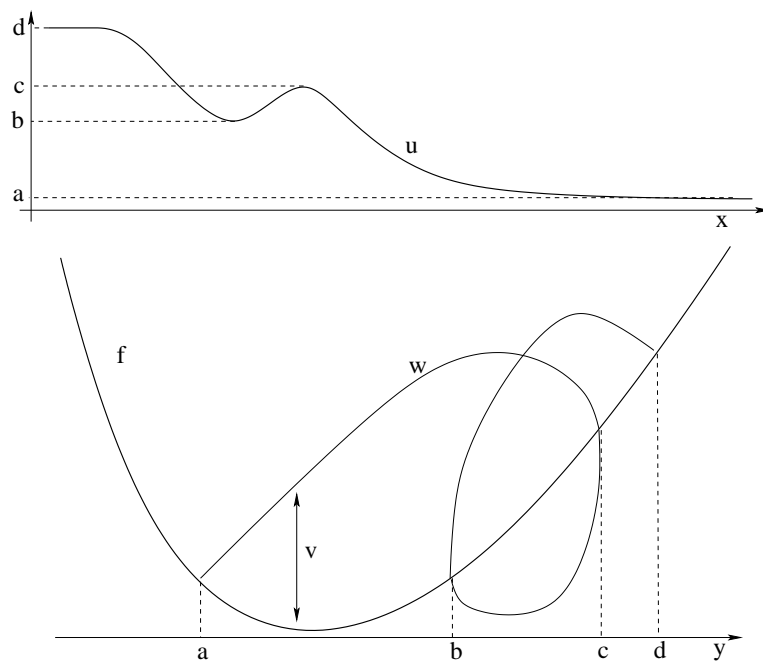
$$\begin{aligned} \int_0^\infty \int_{\mathbb{R}} |z(t, x)| |z^*(t, x)| dx dt &\leq - \int_0^\infty \left[\frac{d}{dt} \frac{Q(t)}{c} \right] dt \\ &\leq \frac{Q(0)}{c} \\ &\leq \frac{1}{c} \int_{\mathbb{R}} |z(0, x)| dx \cdot \int_{\mathbb{R}} |z^*(0, x)| dx \end{aligned}$$

Lyapunov Functionals for the Scalar Viscous Conservation Law

$$u_t + f(u)_x = u_{xx}$$

The curve $\gamma \doteq \begin{pmatrix} u \\ f(u) - u_x \end{pmatrix} = \begin{pmatrix} \text{conserved quantity} \\ \text{flux} \end{pmatrix}$

evolves according to $\gamma_t + f'(\gamma_x) = \gamma_{xx}$



Curve Length

$$\gamma \doteq \begin{pmatrix} u \\ f(u) - u_x \end{pmatrix} \quad \gamma_x = \begin{pmatrix} v \\ w \end{pmatrix} \doteq \begin{pmatrix} u_x \\ -u_t \end{pmatrix}$$

Length $L(\gamma) \doteq \int |\gamma_x| dx = \int \sqrt{v^2 + w^2} dx$

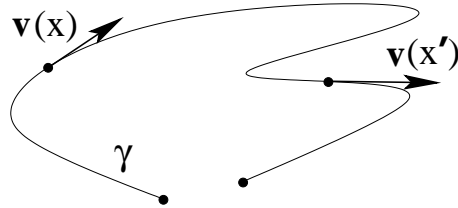
$$\gamma_t + f'(u)\gamma_x = \gamma_{xx} \quad \implies \quad \frac{dL}{dt} \leq 0$$

$$\begin{aligned} -\frac{d}{dt}L(\gamma(t)) &= \int_{\mathbb{R}} \frac{|v| [(w/v)_x]^2}{(1 + (w/v)^2)^{3/2}} dx \\ &\geq \frac{1}{(1 + \delta^2)^{3/2}} \int_{|w/v| \leq \delta} |v| [(w/v)_x]^2 dx \end{aligned}$$

$$\begin{aligned} \int_0^\infty \int_{|w/v| \leq \delta} |v| [(w/v)_x]^2 dx dt &= \mathcal{O}(1) \cdot \int_0^\infty \left| \frac{d}{dt}L(\gamma(t)) \right| dt \\ &= \mathcal{O}(1) \cdot L(\gamma(0)) \end{aligned}$$

Area Swept by a Curve

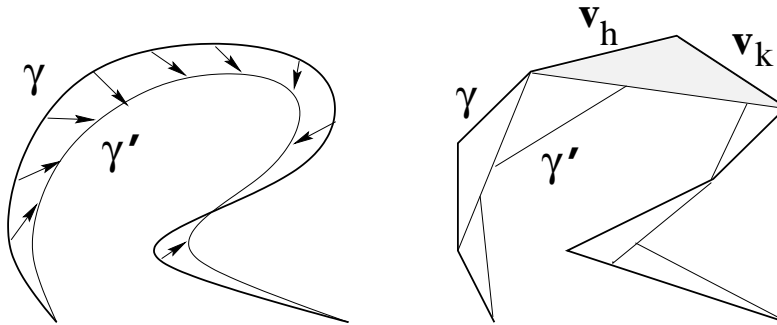
$$\gamma_t + f'(u)\gamma_x = \gamma_{xx}$$



Area functional $Q(\gamma) \doteq \frac{1}{2} \iint_{x < x'} |\gamma_x(x) \wedge \gamma_x(x')| dx dx'$

controls the area swept by the curve γ : $|dA| \leq -dQ$

(for a proof, think of polygonal approximations)



$$Q(\gamma) = \frac{1}{2} \sum_{i < j} |\mathbf{v}_i \wedge \mathbf{v}_j|$$

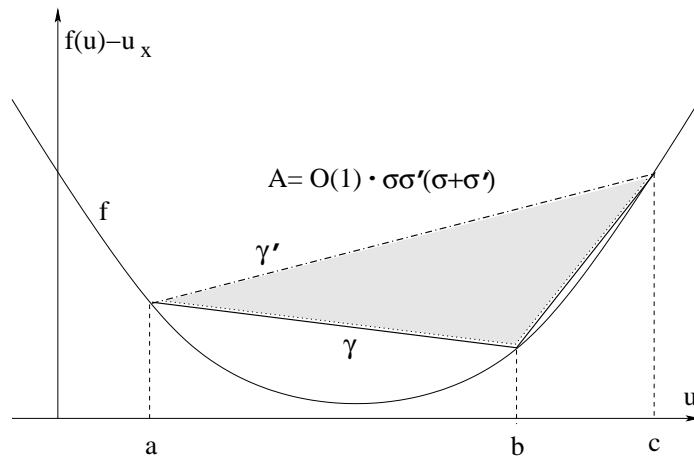
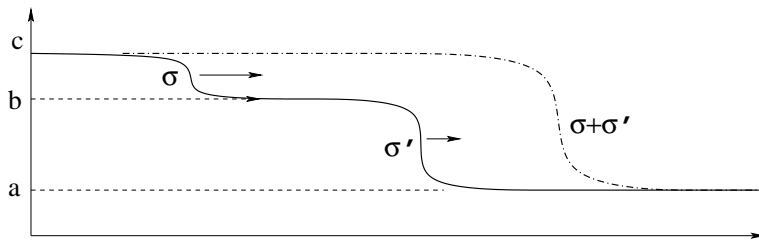
$$|dA| = \frac{1}{2} |\mathbf{v}_h \wedge \mathbf{v}_k| \leq -dQ$$

$$\gamma_t + f'(u)\gamma_x = \gamma_{xx} \quad \gamma_x = \begin{pmatrix} v \\ w \end{pmatrix} \doteq \begin{pmatrix} u_x \\ -u_t \end{pmatrix}$$

$$-\frac{dQ}{dt} \geq \left| \frac{dA}{dt} \right| = \int |\gamma_t \wedge \gamma_x| dx = \int |\gamma_{xx} \wedge \gamma_x| dx = \int |v_x w - v w_x| dx$$

$$\int_0^\infty \int |v_x w - v w_x| dx dt \leq \int_0^\infty \left| \frac{dQ(\gamma(t))}{dt} \right| dt \leq Q(\gamma(0))$$

Genuinely Nonlinear Case



Theorem (S. Bianchini, 2001). Consider a strictly hyperbolic system with viscosity

$$u_t + A(u)u_x = \varepsilon u_{xx} \quad u(0, x) = \bar{u}(x). \quad (10)$$

There exists $\delta > 0$ such that the following holds. If

$$\text{Tot.Var.}\{\bar{u}\} < \delta, \quad \|\bar{u}\|_{\mathbf{L}^\infty} < \delta,$$

then the Cauchy problem (10) admits a unique solution $u^\varepsilon(t, \cdot) = S_t^\varepsilon \bar{u}$, defined for all $t \geq 0$. Moreover, we have the estimates

$$\text{Tot.Var.}\{S_t^\varepsilon \bar{u}\} \leq C \text{Tot.Var.}\{\bar{u}\}, \quad (\mathbf{BV} \text{ bounds})$$

$$\|S_t^\varepsilon \bar{u} - S_t^\varepsilon \bar{v}\|_{\mathbf{L}^1} \leq L \|\bar{u} - \bar{v}\|_{\mathbf{L}^1} \quad (\mathbf{L}^1 \text{ stability})$$

(Convergence) As $\varepsilon \rightarrow 0$, the solutions u^ε converge to the trajectories of a semi-group S , such that

$$\|S_t \bar{u} - S_t \bar{v}\|_{\mathbf{L}^1} \leq L \|\bar{u} - \bar{v}\|_{\mathbf{L}^1} \quad t \geq 0.$$

These vanishing viscosity limits can be regarded as the unique **viscosity solutions** of the hyperbolic Cauchy problem

$$u_t + A(u)u_x = 0 \quad u(0, x) = \bar{u}(x).$$

- In the conservative case $A(u) = Df(u)$, the viscosity solutions are weak solutions to

$$u_t + f(u)_x = 0 \quad u(0, x) = \bar{u}(x)$$

satisfying the Liu entropy conditions.

- In the case where each field is genuinely nonlinear or linearly degenerate, the viscosity solutions coincide with the unique limits of Glimm or front tracking approximations.

Concluding Remarks

- Alternative approaches to systems of conservation laws:
 1. The BV setting (Glimm). Control of the total variation achieved by an interaction potential.
 2. The L^p setting (DiPerna). Convergence of approximate solutions achieved by compensated compactness.

- In the BV setting the Cauchy problem is well posed. Examples show it can be ill posed for L^∞ data.

- Convergence of vanishing viscosity solutions can also be achieved in BV , controlling the total variation by suitable Lyapunov functionals.

- Riemann Problem (and interactions of elementary waves) is no longer at the center stage.

- Greater generality (no assumptions about genuine nonlinearity).

- Reduce the gap between analytical theory and numerical algorithms.